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OSCILLATION OF CONDUCTION ELECTRON DENSITY  
NEAR THE SOLUTE ATOMS IN DILUTE Cu-Mn ALLOY

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OSCILLATION OF CONDUCTION ELECTRON DENSITY  
NEAR THE SOLUTE ATOMS IN DILUTE Cu-Mn  
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The "wipe-out" number characterizing the first order quadrupole effect of the  $^{63}\text{Cu}$  NMR spectrum in dilute Cu-Mn alloys is  $1500 \pm 75$ . This value agrees quite well with the one calculated in terms of the Blandin-Caroli<sub>A</sub> resonance scattering theory.

The properties of transition element impurities in nonmagnetic hosts have been the subject of numerous investigations [1], although the level of understanding is still unsatisfactory. This paper seeks to demonstrate that the measurement of the asymptotic excess conduction electron density /Friedel oscillation/ contributes to the solution of the problem of the "magnetic" impurity in a nonmagnetic host. The aim has been to observe Friedel oscillation around "magnetic" solute Mn atoms in copper-based alloys by measuring the first order quadrupole effect appearing in the  $^{63}\text{Cu}$  NMR spectrum [2], and preliminary results for this Cu-Mn system are presented here. In this case the effect of Friedel oscillation is easily separated from that of the spin density oscillation. The theoretical prediction [3] that the Friedel oscillation is nonzero has been experimentally confirmed by the second order quadrupole measurements of Lumpkin [4].

The master alloy was prepared from 99,999 atomic % copper and 99,9 atomic % manganese by vacuum smelting. The measurements were performed on  $\sim 15\mu$ -thick foils produced by cold rolling. To remove dislocation effects the specimens were annealed for 1 hour at  $400^\circ\text{C}$ . The short duration and the relatively low temperature of the annealing were chosen to avoid the internal oxidation observed by Howling [5]. The Mn concentrations of the specimens prepared from the master alloy, as determined by spectro-

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\* The paper will appear in Low Temperature Physics, LT-12, Kyoto 1970.

photometric analysis, were found to be 52, 142, 270, 460, and 1040 atomic ppm /  $\pm 5\%$  /. The wide line NMR spectrometer described in [6] was used.

The wipe-out number was determined from the NMR amplitude /D/-impurity concentration /c/ curve /Fig. 1/. The experimental points are averages taken over about 50 spectra. Simultaneously measured  $^{27}\text{Al}$  NMR signals from Al foils were used as reference. The measured D amplitudes have been normalized to that of pure copper /  $D_0$  /, and the satellite contribution to the peak - to - peak amplitude of the derivative signal is given. The results measured at  $T = 150^\circ\text{C}$  are shown. Below this temperature the  $\log \left( \frac{D}{D_0} - 0,4 \right) / 0,6 - c$  curve deviates from the straight line characteristic of the first order quadrupole effect [7] owing to the spin density oscillation.

At  $150^\circ\text{C}$  our experiments gave the wipe-out number  $n_1 = 1500 \pm 75^*$ .

Assuming the scattering to be predominantly d-resonance scattering, the excess density of the conduction electrons with spin near the impurity is given by [3] as

$$\Delta\rho^\sigma(r) = -\frac{5}{4\pi^2} \sin\delta_2^\sigma r^{-3} \cos(2k_F r + \delta_2^\sigma) ,$$

where  $r$  is the separation of matrix nucleus from impurity,  $k_F$  the Fermi wave number. The phase shifts  $\delta_2^\sigma$  of the scattered partial waves  $l = 2$ , taken at the Fermi energy, are related to the number of localized 3d-electrons by the expression  $N = \frac{5}{\pi} (\delta_2^\uparrow + \delta_2^\downarrow)$ , and to the  $z$  component of the magnetic moment localized on the impurity atoms by the formula  $M_z = \frac{5\mu_B}{\pi} (\delta_2^\uparrow - \delta_2^\downarrow)$ ,  $\mu_B$  is the Bohr magneton./

The difference in the number of spin-up and spin-down electrons gives the spin density, while their sum gives the charge density oscillation. Thus, the component  $z$  of the electric field gradient can be expressed as

$$q = \frac{8\pi}{3} \alpha [\Delta\rho^\uparrow(r) + \Delta\rho^\downarrow(r)]$$

where  $\alpha = 25$  is the enhancement factor [8, 9].

$$\delta_2^\uparrow = \pi; \quad \delta_2^\downarrow = \frac{\pi}{5} \cdot 0,7 = 0,44$$

Using the phase shifts given by Daniel [10] in his calculation of the residual resistivity in Cu-3d alloys, and the computer-calculated result [7] which relates the Friedel oscillation directly to the wipe-out number, we obtain the theoretical value  $n_1 \text{ theor} = 1420$ . This is in good agreement with the experimental value. /In terms of the "all-or-nothing" language the  $q_{\text{krit}} = 1,8 \times 10^{21} \text{ cm}^{-3}$  is the same as for Cu-Zn [7]./

The electronic structure proposed by Hurd [11] gives  $n_1 \text{ theor} = 1840$

\* The extrapolation method and the wipe out number published in the KFKI preprint 9/1969 are incorrect.

in satisfactory agreement with the experimental value.

The quadrupole effect of the NMR of the matrix nuclei in dilute alloys of copper with 3d - transition metals can be easily estimated from resonance scattering theory, e.g. by using the phase shifts given in [10] or by making use of assumption about the electronic structure of 3d - transition metals embedded in copper [1, 12].

Above the Kondo temperature the dependence of the excess charge density near the solute atoms on the atomic number of the latter yields a v-shaped curve with manganese at the minimum.

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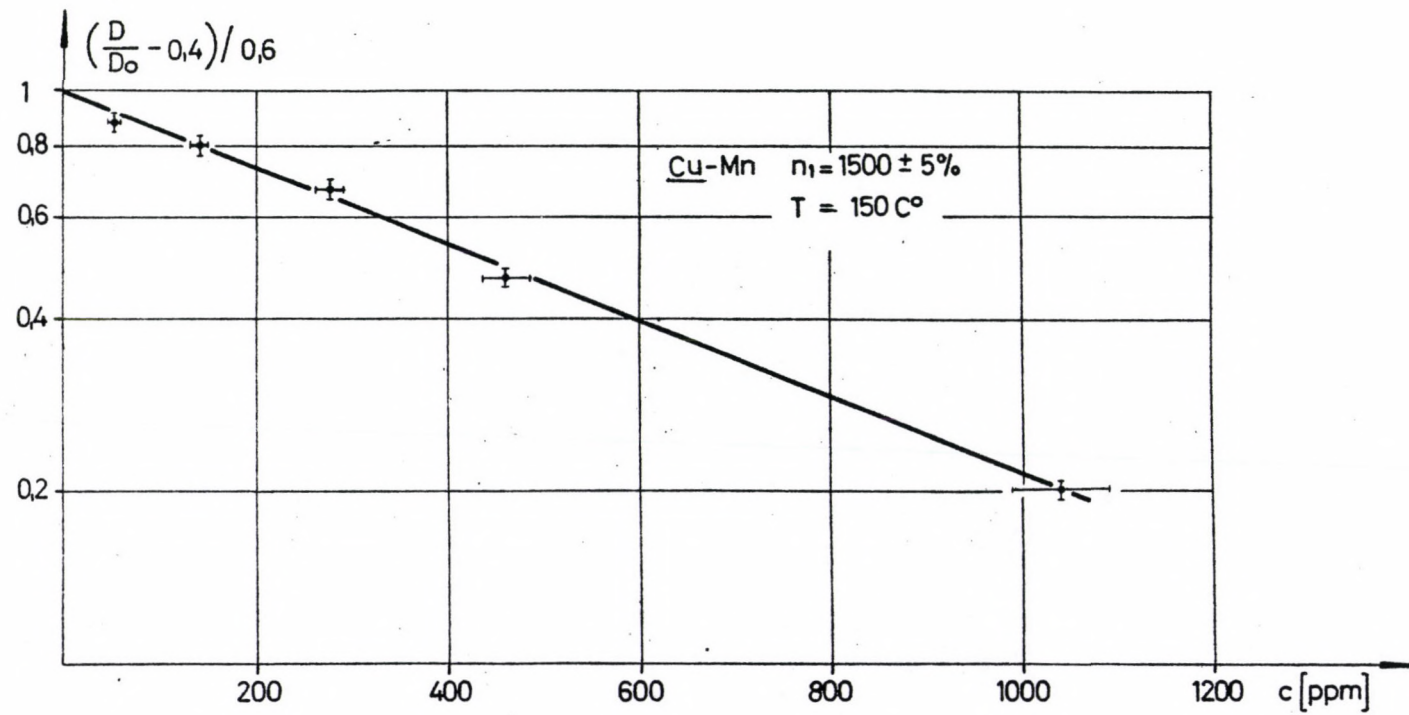
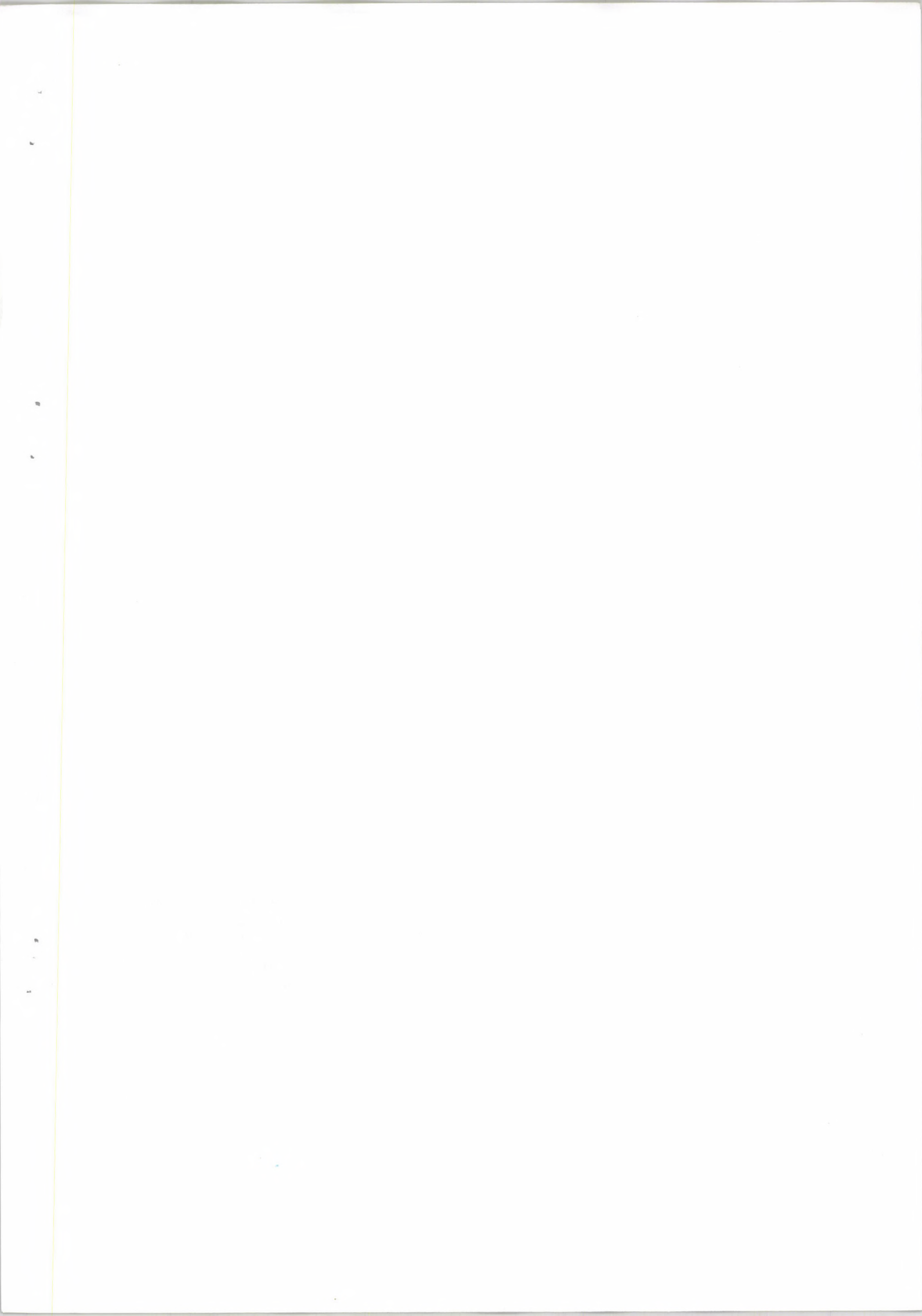


Figure 1.

First order quadrupole effect in copper based Cu-Mn dilute alloys, at  $T = 150\text{ C}^\circ$



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